

NUMERICAL STUDY OF TUBE-BUNDLE FLOW-INDUCED VIBRATIONS WITH MULTIPHASE-POD APPROACH

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Abstract. *Fluid-Structure Interactions are present in a large number of systems of nuclear power plants and nuclear on-board stoke-holds. Particularly in steam generators, where tube bundles are submitted to cross-flow which can lead to structure vibrations. We know that numerical studies of such a complex mechanism is very costly, that is why we propose the use of reduced-order methods in order to reduce calculation times and to make easier parametric studies for such problems. We use the multiphase-POD approach, which is an adaptation of the classical POD approach to the case of a moving structure in a flow, considering the whole system (fluid and structure) as a multiphase domain. We are interested in the case of large displacements of a structure moving in a fluid, in order to observe the ability of the multiphase-POD technique to give a satisfying solution reconstruction. We obtain very interesting results for the case of a single circular cylinder in cross-flow (lock-in phenomenon). Then we present the application of the method to a case of confined cylinders in large displacements too. Here again, results are encouraging. An on-going work consist in going further testing parametric studies with POD-Galerkin approach and with POD basis interpolation. A future work will consist in applications to fluid-structure interactions.*

Keywords: *Tube bundle systems, Multiphase-POD, Flow-Induced Vibrations.*

1. INTRODUCTION

Nuclear power plants and nuclear on-board stokeholds are containing very complex installations where flow-induced vibrations are present at various levels. Particularly, tube bundle systems in the steam generator part are prone to fretting-wear or even breaking, mainly due to transverse flow-induced vibration problems [4, 7, 8, 13, 18, 25, 23, 24]. But, as well as experiments, numerical simulations of fluid-structure interactions in a tube bundle system remain very costly [5, 15, 28] : in order to give a relevant description of the flow and of the structure displacements, it is necessary to construct a system with a large number of degrees of freedom. The three-dimensional aspect of the flow and of the fluid-structure interactions

added to the large number of tubes and their reciprocal interactions force us to lead long-time calculations, in order to be as precise as possible.

One of the existing alternatives to these restrictive situations from an industrial point of view is to have recourse to reduced-order models [1, 11, 2, 9, 19, 21] such as POD (Proper Orthogonal Decomposition) [29, 14, 17]. It is thus possible to make numerical calculations in a very short time, and this paves the way to lead parametric studies or even real-time control. Here, the challenge consists in the adaptation of POD to the case of fluid-structure interaction problems. Liberge & Hamdouni [19, 20] proposed an efficient way to cope with the moving fluid-structure interface; this method is called ‘‘Multiphase-POD’’ : the idea is to consider the whole system (fluid and structure domains) as a unique multiphase domain. The advantages of this technique are numerous. First, the data can be extracted from any way (experiments, moving grid techniques, etc.), it is only necessary to know how data are organized to proceed to the interpolation. Here, we show its ability to reproduce large structure displacements, which is often, to the author’s knowledge, not so easy with other reduced-order models. The first part of this paper is dedicated to the description of the Multiphase-POD method. Then, its application to the case of large displacements of a single circular cylinder under cross-flow (lock-in phenomenon) is presented in a second part. Finally, a 2D tube-bundle configuration is considered with one moving cylinder under cross-flow.

2. MULTIPHASE-POD

We consider here that POD-Galerkin method is well known (see for example [14]) and we are interested on its adaptation to FSI through Multiphase-POD. Complete calculations are leaded with a classic ALE approach [12]. Thus, in this case of Flow Induced Vibrations, classic POD-Galerkin method cannot be used because of the presence of a moving interface : POD modes are only spatial and consequently, they do not contain any dynamic information, although snapshots, in the case of an ALE calculation, have been taken for several positions of mesh nodes.

To get round this problem, Liberge & Hamdouni [19] proposed an original method that treats the case of a fluid-structure interaction problem with an adaptation of the POD-Galerkin technique, which is called ‘‘Multiphase-POD method’’, where a non-moving mesh is used. The description of the Multiphase-POD method is the following : lets consider a global domain Ω containing the fluid domain $\Omega_f(t)$ and the solid domain $\Omega_s(t)$ at each time step t , where the solid domain is considered as a particular fluid with its own physical characteristics (density, viscosity). We have $\Omega = \Omega_f(t) \cup \Omega_s(t) \cup \Gamma_i(t)$, where $\Gamma_i(t)$ is the interface between fluid and solid domains. A global velocity field $u \in H(\Omega)$ (with H a Hilbert space) is considered :

$$u(x, t) = u_f(x, t)\chi_{\Omega_f}(x, t) + u_s(x, t)\chi_{\Omega_s}(x, t) \quad (1)$$

where χ_{Ω_f} and χ_{Ω_s} are respectively characteristics functions defining if the considered point position is in the fluid or in the solid domain :

$$\chi_{\Omega_s}(x, t) = \begin{cases} 1 & \text{if } x \in \overline{\Omega_s} \\ 0 & \text{if } x \notin \overline{\Omega_s} \end{cases} \quad (2)$$

and $\chi_{\Omega_f}(x, t) = 1 - \chi_{\Omega_s}(x, t)$. Taking into account this notations, a global weak form of Navier-Stokes equations on Ω is made possible to formulate :

$$\int_{\Omega} \rho \frac{\partial u(x, t)}{\partial t} u^* dx + \int_{\Omega} (u \cdot \nabla) u \cdot u^* dx = \int_{\Omega} (\nabla \cdot \sigma) u^* dx \quad (3)$$

where u^* is a test-function defined as $u^* \in H(\Omega)$ with the non-deformable solid constraint :

$$D(u^*) = 0 \text{ in } \Omega_s(t) \quad (4)$$

Each field or variable is defined on the global domain Ω as described below:

$$\begin{cases} u(x, t) = u_f(x, t)\chi_{\Omega_f}(x, t) + u_s(x, t)\chi_{\Omega_s}(x, t) \\ \sigma(x, t) = \sigma_f(x, t)\chi_{\Omega_f}(x, t) + \sigma_s(x, t)\chi_{\Omega_s}(x, t) \\ \rho(x, t) = \rho_f(x, t)\chi_{\Omega_f}(x, t) + \rho_s(x, t)\chi_{\Omega_s}(x, t) \\ \mu(x, t) = \mu_f(x, t)\chi_{\Omega_f}(x, t) + \mu_s(x, t)\chi_{\Omega_s}(x, t) \end{cases} \quad (5)$$

Lets define both components of the constraints tensor σ :

$$\sigma_{f,ij}(x, t) = -p\delta_i^j + 2\mu_f D_{ij}(u_f) \quad (6)$$

where δ_i^j is the Kroenecker symbol and D_{ij} is the deformation velocity tensor. The definition of the structural compotent $\sigma_s(x, t)$ allows taking into account that the solid has its specific viscosity and the non-deformable structural condition. For the viscosity, a penalization term is used: in order to specify that the domain $\Omega_s(t)$ is solid, the viscosity is artificially increased. To insure the non-deformable condition, a Lagrange multiplier Λ is added. Thus, the structural component of the constraints tensor is:

$$\sigma_{s,ij}(x, t) = -p\delta_i^j + \Lambda + 2\mu_s D_{ij}(u_s) \quad (7)$$

Developing the global weak form with these definitions and making the Proper Orthogonal Decomposition on the global velocity flow field leads to the construction of a dynamical system for the whole domain Ω which is fixed all over the studied time interval. Taking into account the space-time decomposition of the global velocity field as:

$$u(x, t) = \sum_{i=1}^N a_n(t)\Phi_n(x) \quad (8)$$

where Φ_n , $n = 1, \dots, N$ are elements of the POD basis and $a_n(t)$, $n = 1, \dots, N$ are time coefficients, the final dynamical system is the following :

$$\begin{cases} \sum_{i=1}^N \frac{da_i}{dt} A_{in} = \sum_{i=1}^N \sum_{j=1}^N B_{ijn} a_i a_j + \sum_{i=1}^N C_{in} a_i + E_n \\ D(u) = 0 \text{ on } \Omega_s(t) \quad \text{(non deformability)} \\ \frac{\partial \chi_{\Omega_s}}{\partial t} + u \cdot \nabla \chi_{\Omega_s} = 0 \quad \text{(characteristic function transfer)} \end{cases} \quad (9)$$

for each $n = 1, \dots, N$ where N is the number of modes in the POD basis. Coefficients A_{in} , B_{ijn} , C_{in} , E_n are not detailed here, but a very important point to notice is that they are not

all exclusively spatial coefficients, because some of them contain the physical characteristics $\rho(x, t)$ and $\mu(x, t)$. Thus, they have to be re-calculated at each time step: the time calculation is increased in comparison with a classic POD model without moving structure. But this time calculation is still less than a complete calculation. Another approach consists in making the proper orthogonal decomposition of the characteristic function $\chi_{\Omega_f}(x, t)$ also, which allows avoiding the time dependence of all coefficients of the dynamical system. For more precisions, see [20]. Practical implementation of the Multiphase-POD technique is described below.

1. Lead a complete ALE calculation of the fluid-structure interaction problem during a time interval $[0, T]$
2. Extract enough snapshots from this complete calculation
3. Create a unique Cartesian fixed mesh containing both fluid and solid domains
4. Interpolate each extracted snapshot onto the fixed reference mesh: new fixed snapshots are created
5. Apply the classic POD approach for the new snapshots constructed on the reference mesh
6. Construct the dynamical system following (9) and resolve it with a classic method (Runge-Kutta for example).

3. APPLICATION TO LOCK-IN PHENOMENON OF A SINGLE CIRCULAR CYLINDER

In a first time, we propose a simple application of the Multiphase-POD technique, which is the case of a circular cylinder under cross-flow (see Fig.1). The cylinder is submitted to transverse displacements (y -direction) due to the presence of the flowing fluid. The case of small displacements of the structure has already been tested with Multiphase-POD in [19]. Here, we consider the lock-in phenomenon [16, 26, 30], where amplitudes of the structure displacement are of the order of the cylinder radius. The fluid domain is considered as infinite, as boundaries are far enough from the structure. The effects that the flow exerts on the structure are modeled through a restoring force. Reynolds number is $\mathcal{R}_e = 100$, fluid is water. Cylinder displacement maximal amplitude is $A^* = 0.58D$, where D is the cylinder diameter: the frequency lock-in mechanism is reached.

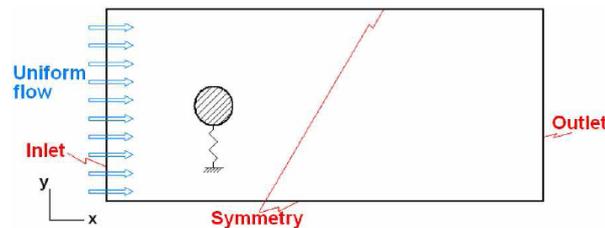


Figure 1. Case of a single circular cylinder and boundary conditions

Complete calculations are leaded with the CFD code *Code_Saturne* [3] and data at the interface are interpolated to the cylinder gravity center. The reduced-order model is constructed with the following characteristics: 250 snapshots are extracted from the complete ALE calculation, 6 POD modes are constituting the POD basis. The fixed reference mesh contains 200 x 250 points. The dynamical system resolution in the present case is simplified: indeed, the penalization term is sufficient to guarantee the non-deformable condition. Time integration scheme is Runge-Kutta 4.

The two first time coefficients are represented on Fig.2, they are well reconstructed by the reduced model. And, as they are containing the main part of the system energy, this good reproduction allows a good reconstruction of the velocity flow field and the cylinder displacement is also well reproduced (Fig.3), which is confirming that 1) the Multiphase-POD method is able to reproduce a structure displacement and a fluid flow with its global formulation and 2) the Multiphase-POD method is able to reproduce large displacements of the structure. The latter point is interesting for the willingness of studying instability behaviors.

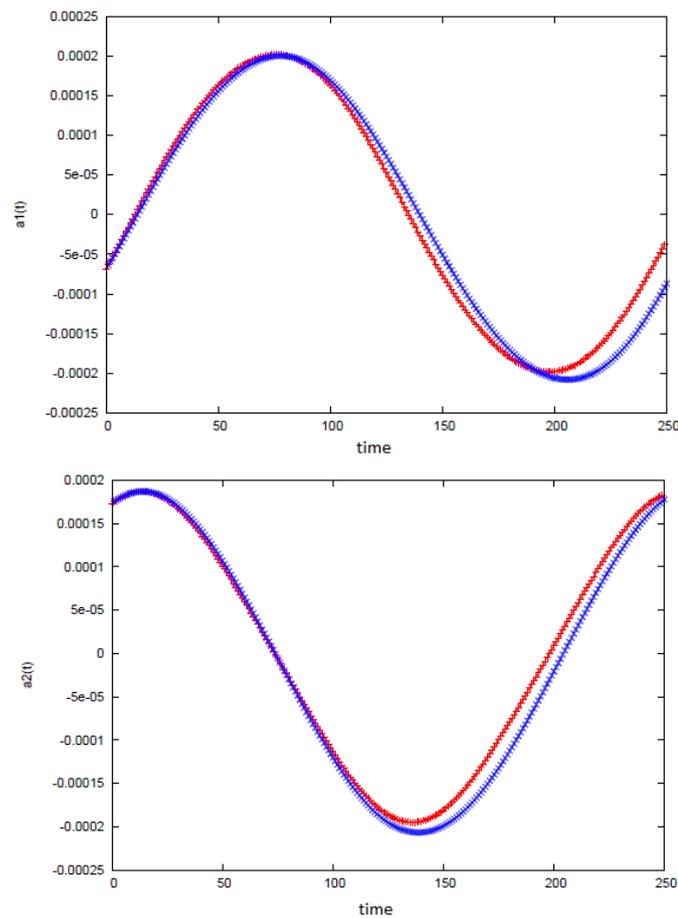


Figure 2. Two first time coefficients of the velocity field for the single cylinder. +++ direct coefficient; xxx Multiphase-POD reconstruction

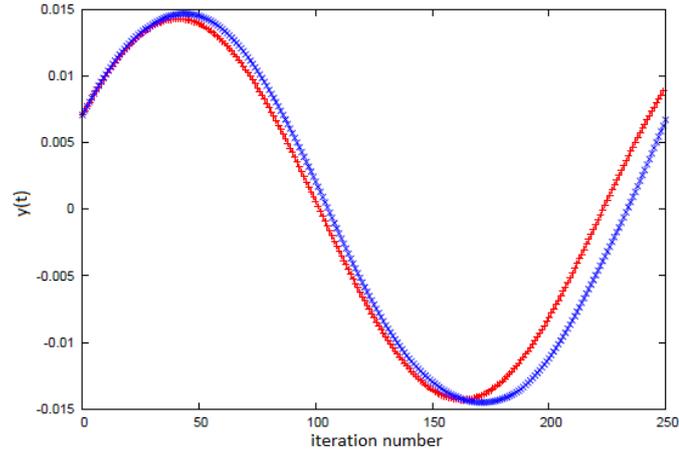


Figure 3. Gravity center displacement reconstruction of the cylinder. +++ direct calculation; xxx Multiphase-POD reconstruction

4. APPLICATION TO FIV IN TUBE-BUNDLE CONFIGURATION

In order to consider a configuration close to the case of a tube bundle of heat exchanger, we consider a circular cylinder in a confined configuration. Non-dimensional numbers are adapted to this configuration, here Reynolds number is defined as $\mathbf{R}_e = \frac{\rho U_p D}{\mu}$. The step fluid velocity U_p takes into account the tube confinement and is defined as: $U_p = U_\infty \frac{P}{P - D}$ where U_∞ is the equivalent mean flow velocity that would have been imposed in an infinite domain and P is the pitch ratio (distance between two neighbouring cylinders centres). Geometry and boundary conditions are depicted on Fig.4: a 2D domain and only one tube and its neighbors are considered, with inlet/outlet boundary conditions. Thus, the domain is not representing a whole tube bundle but a confined case. Reynolds number is fixed to $\mathcal{R}_e = 2000$, complete calculation is also leaded with *Code_Saturne* which has been validated in various FSI studies in tube bundle systems [5, 15, 22]. Large displacements in the y -direction (see Fig.4) of the central cylinder are considered ($A^* = 0.35D$ when $P/D = 0.44D$). The reference fixed mesh contains 200 x 200 points.

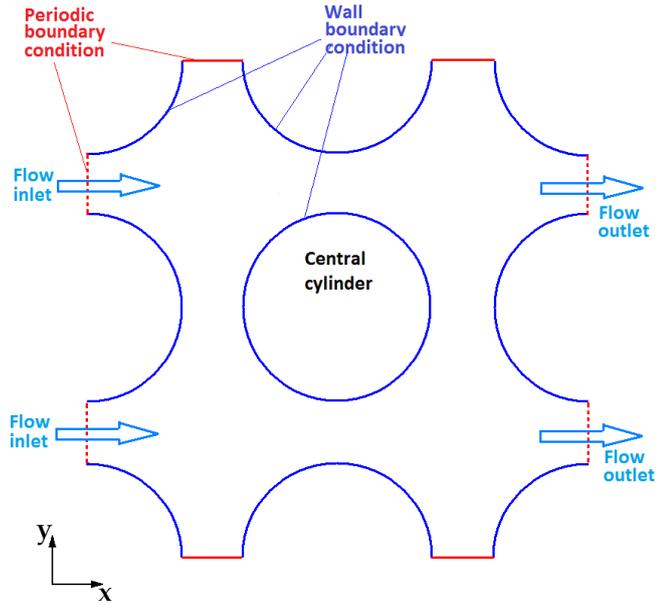


Figure 4. Boundary conditions for the confined tube

Figure 5 represents the comparison between the global velocity flow field from the complete calculation and the interpolated velocity flow field. It allows to check the precision of the snapshots interpolation algorithm: velocity levels are well reproduced after interpolation. In the interpolated case (right-hand side), the non-zero velocity in the cylinder zone is representative of the structure velocity, which is now considered as the second phase of the flow.

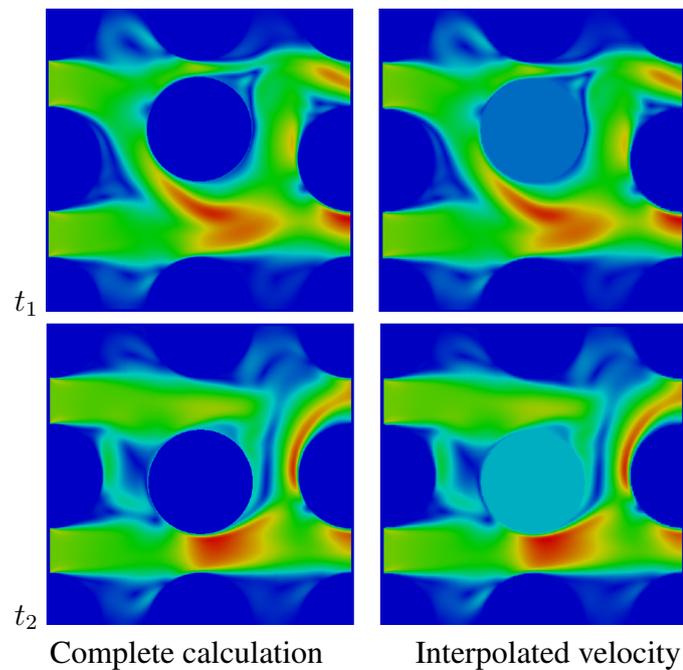


Figure 5. Comparison between complete and interpolated velocity field at two dates t_1 and t_2

Figure 6 shows the comparison between the central cylinder displacement calculated by complete calculation and by Multiphase-POD. The reconstruction gives very satisfying results, which is confirmed by the observation of the two first time coefficients of the global velocity flow field (Fig.7). The reconstruction of large displacements with Multiphase-POD in the case of a confined tube bundle is very interesting. It allows to plan for its implementation to unstable fluid-structure interactions like fluid-elastic instability occurring in tube bundle systems [6, 10, 18, 25, 27].

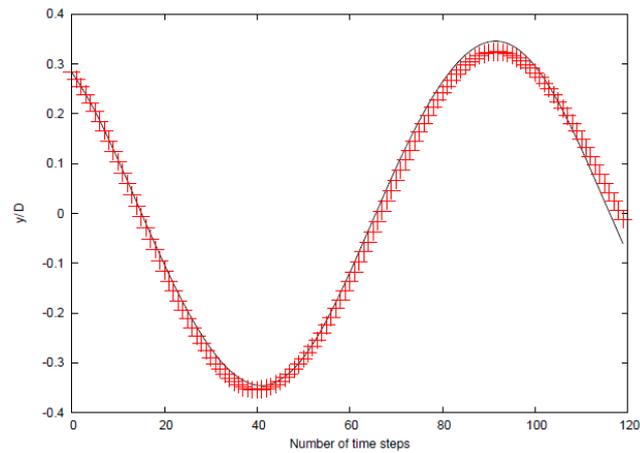


Figure 6. Comparison between complete and multiphase-POD reconstruction of the central cylinder displacement. — complete calculation; +++ Multiphase-POD reconstruction

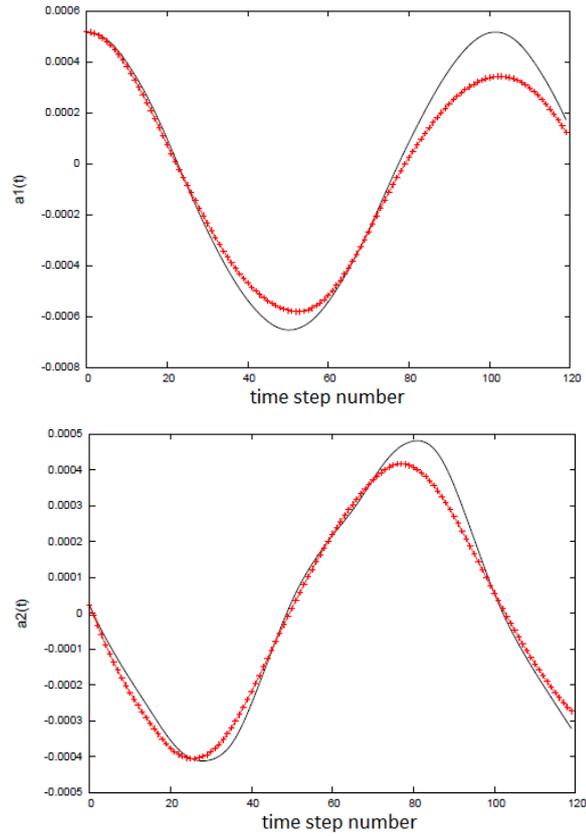


Figure 7. Comparison between complete and multiphase-POD reconstruction of time coefficients. +++ direct coefficient; — Multiphase-POD reconstruction

CONCLUSION

In this paper, the Multiphase-POD method is presented and applied to the case of a single circular cylinder moving under cross-flow and a confined cylinder in tube bundle under cross-flow. The method was already shown to be efficient in the case of small displacements of a structure under flow solicitations and here, we show its efficiency in the case of large displacements of the structure. This is a very interesting point in order to treat instabilities that can appear in a large number of industrial systems. Moreover, a simple case of tube bundle has been successfully past with Multiphase-POD, which is encouraging for reducing calculation times in this context. An on-going work is the application of parametric studies with the help of POD: the main interest of reduced-order models consists in their ability to reconstruct various solutions of a system where one or several parameters have been changed. Indeed, the reconstruction of a solution for which we already have the complete calculation is not sufficient. Further work will consist in the application of these methods to the case of FSI.

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